

Spiral structure and gravitational instabilities in protostellar discs

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27 June 2013 - Beijing - Lin-Shu Symposium

Gravitational instabilities in protostellar discs

- ❖ Conditions for instability
- ❖ Dynamics of self-gravitating discs:
 - ❖ Self-regulation
 - ❖ Local vs global behaviour
- ❖ Numerical uncertainties - convergence, fragmentation
- ❖ Observations of density waves in protostellar discs

Gaseous vs stellar discs

- ❖ Protostellar discs are mostly gaseous
- ❖ Potentially significant differences with respect to the case where the structure is determined mostly by a dissipationless component
 - ❖ Role of Lindblad resonances?
 - ❖ Shocks in the gas might play a significant role (gas is not merely a “passive” component that responds to the structure)

Linear stability criterion

- ❖ Linear dispersion relation

$$(\omega - m\Omega)^2 = c_s^2 k^2 - 2\pi G \Sigma |k| + \kappa^2$$

- ❖ Well known axisymmetric instability criterion:

$$Q = \frac{c_s \kappa}{\pi G \Sigma} < \bar{Q} \approx 1$$

- ❖ Equivalent form of the instability criterion

$$\frac{M_{\text{disc}}(R)}{M_\star} \gtrsim \frac{H}{R}$$

- ❖ Need the disc to be cold and / or massive
- ❖ What are the masses and aspect ratio in actual protostellar discs?

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Are protostellar discs linearly unstable?

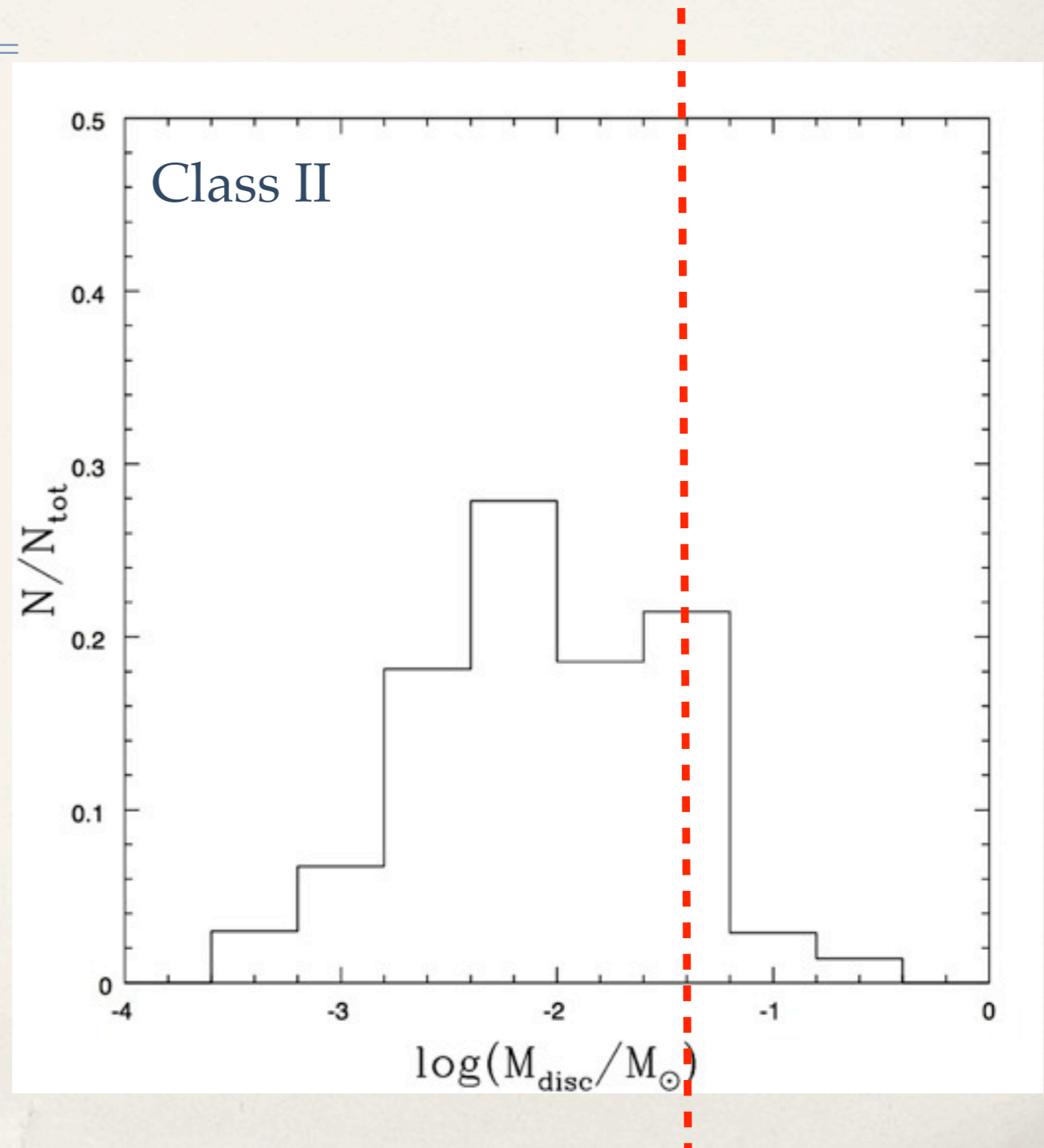
- ❖ Midplane temperature for irradiated discs (Chiang & Goldreich 1997, Chiang & Youdin 2009) gives:

$$\frac{H}{R} \simeq 0.02 \left(\frac{R}{\text{AU}} \right)^{2/7}$$

- ❖ Therefore H/R varies from **0.02** at 1AU to **0.06** at 100 AU
- ❖ Need disc masses of order 5% of the stellar mass to be unstable
- ❖ Protostellar disc masses difficult to measure (see Hartmann et al 2006)

Are protostellar discs linearly unstable?

- * Disc masses in Taurus and Ophiucus by Andrews and Williams (2005, 2007)
- * Disc masses might be underestimated significantly (Hartmann et al 2006)
- * Uncertainties in dust opacities
- * If density profile steep, most of the mass might be hidden in optically thick inner parts (Hartmann 2009)



Are protostellar discs linearly unstable?

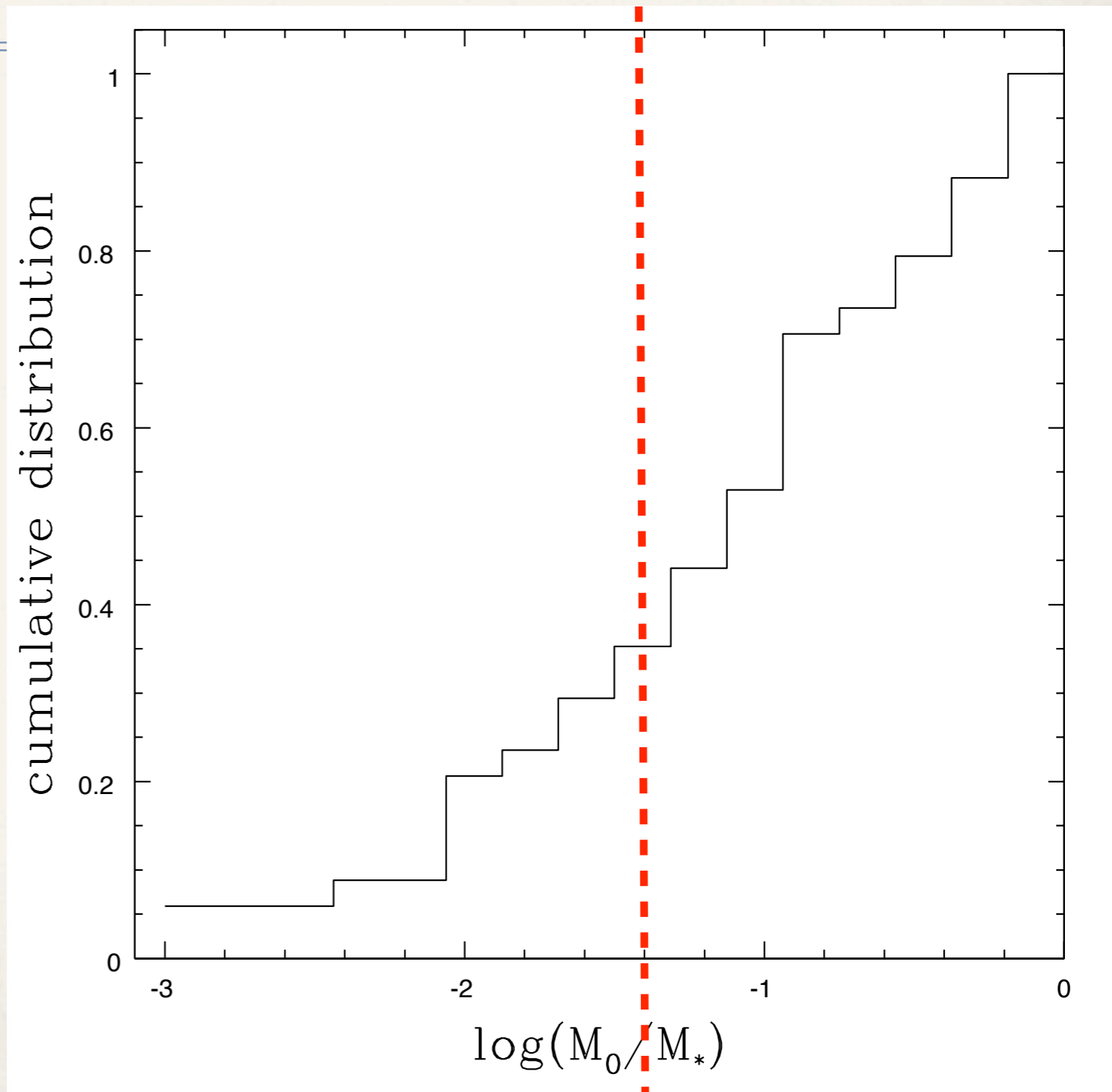
- ❖ Class II (T Tauri) discs are relatively evolved. Can we infer the masses at early stages?
- ❖ Simple (simplistic?) approach:
- ❖ Take all objects with measured M and \dot{M}
- ❖ Apply similarity solutions (Lynden-Bell & Pringle 1973)
- ❖ Find “initial” disc mass and evolutionary timescale

$$M_0 = M_d(t) \left(\frac{t_d}{t_d - t} \right)^{1/2(2-\gamma)}$$

$$t_d = \frac{M_d(t)}{2(2-\gamma)\dot{M}(t)}$$

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 - ❖ Self-regulation
 - ❖ Local vs global behaviour
- ❖ Numerical uncertainties - convergence, fragmentation
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Non linear evolution of GI

- ❖ Investigated numerically in the last decade by several authors (Laughlin & Bodenheimer 1994, Laughlin et al 1998, Pickett et al 2000, Boss 2000, Gammie 2001, Mayer et al 2002, Lodato & Rice 2004, 2005, Mejia et al 2005, Boley et al 2006)
- ❖ Early simulations used an isothermal or polytropic equation of state (Laughlin & Bodenheimer 1994, Mayer et al 2002)
- ❖ Starting from Gammie (2001) it has become clear that the evolution is strongly dependent on the cooling time t_{cool}
- ❖ Introduce a cooling parameter as the ratio of cooling to dynamical timescale

$$\beta = t_{cool}\Omega$$

Thermal self-regulation of GI

- ❖ Role of cooling time clear if one thinks at the form of the stability parameter

$$Q = \frac{c_s \kappa}{\pi G \Sigma} \propto T^{1/2}$$

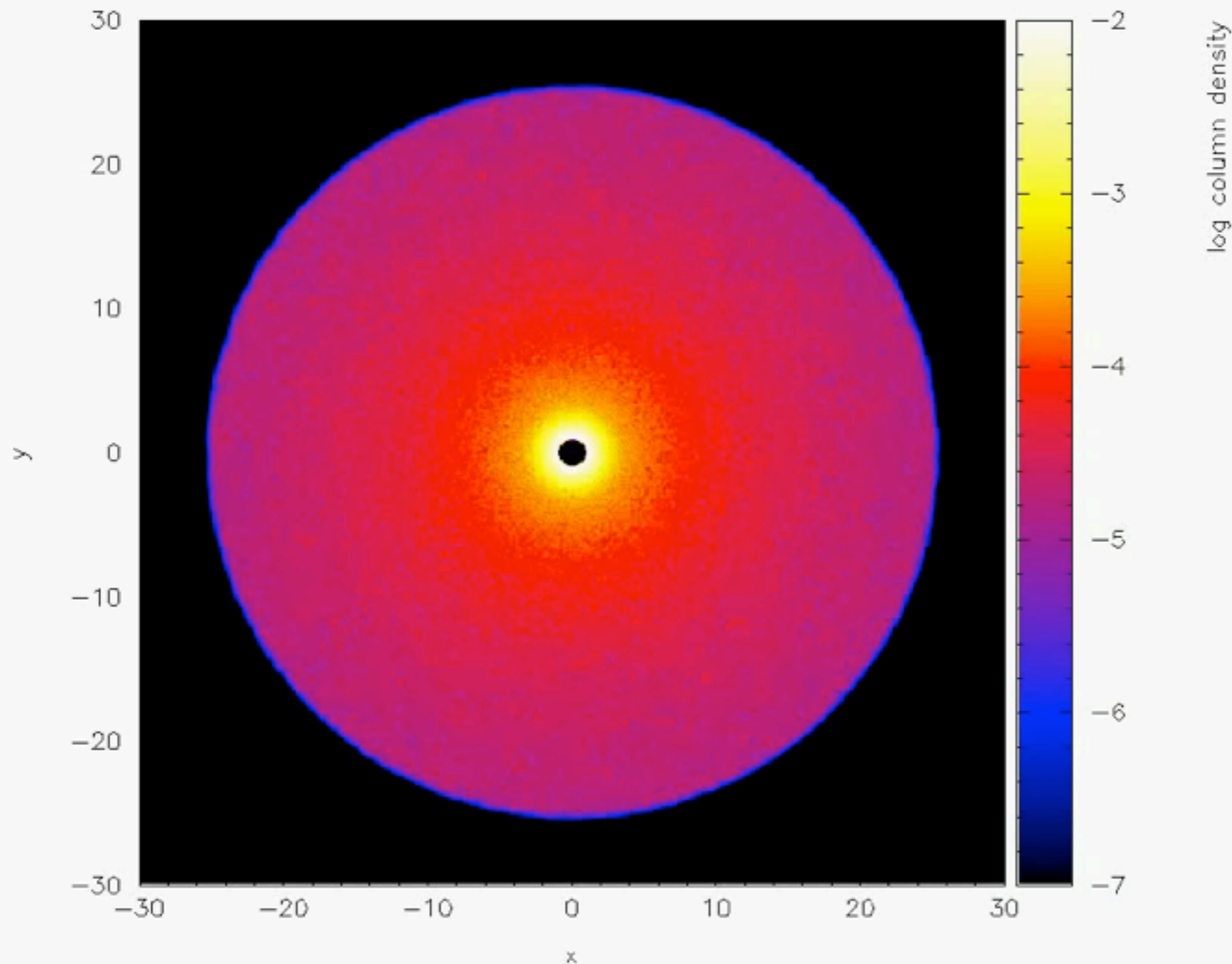
- ❖ Development of the instability feeds energy back onto the equilibrium and stabilizes the disc
- ❖ Works as an effective thermostat for the disc
- ❖ Expect the disc to stay close to marginal stability $Q \sim 1$ (*Paczynski 1977*)
- ❖ Self-regulated discs models can be constructed (*Bertin 1997, Bertin & Lodato 1999*)

Long cooling time: self-regulation

Cossins, Lodato &
Clarke (2009)

$$\beta = 6$$

Long cooling time: self-regulation



Cossins, Lodato &
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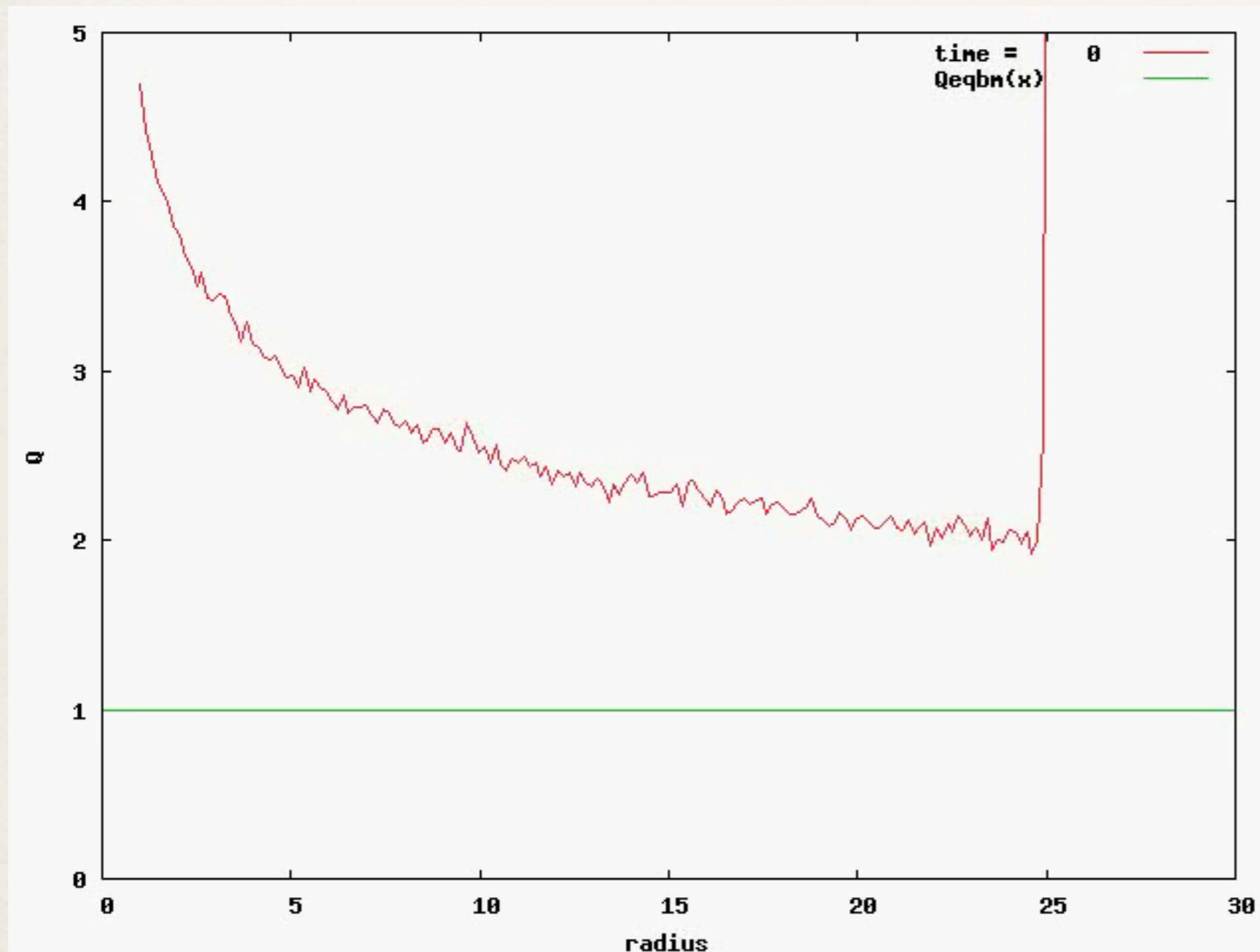
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Thermal saturation of GI

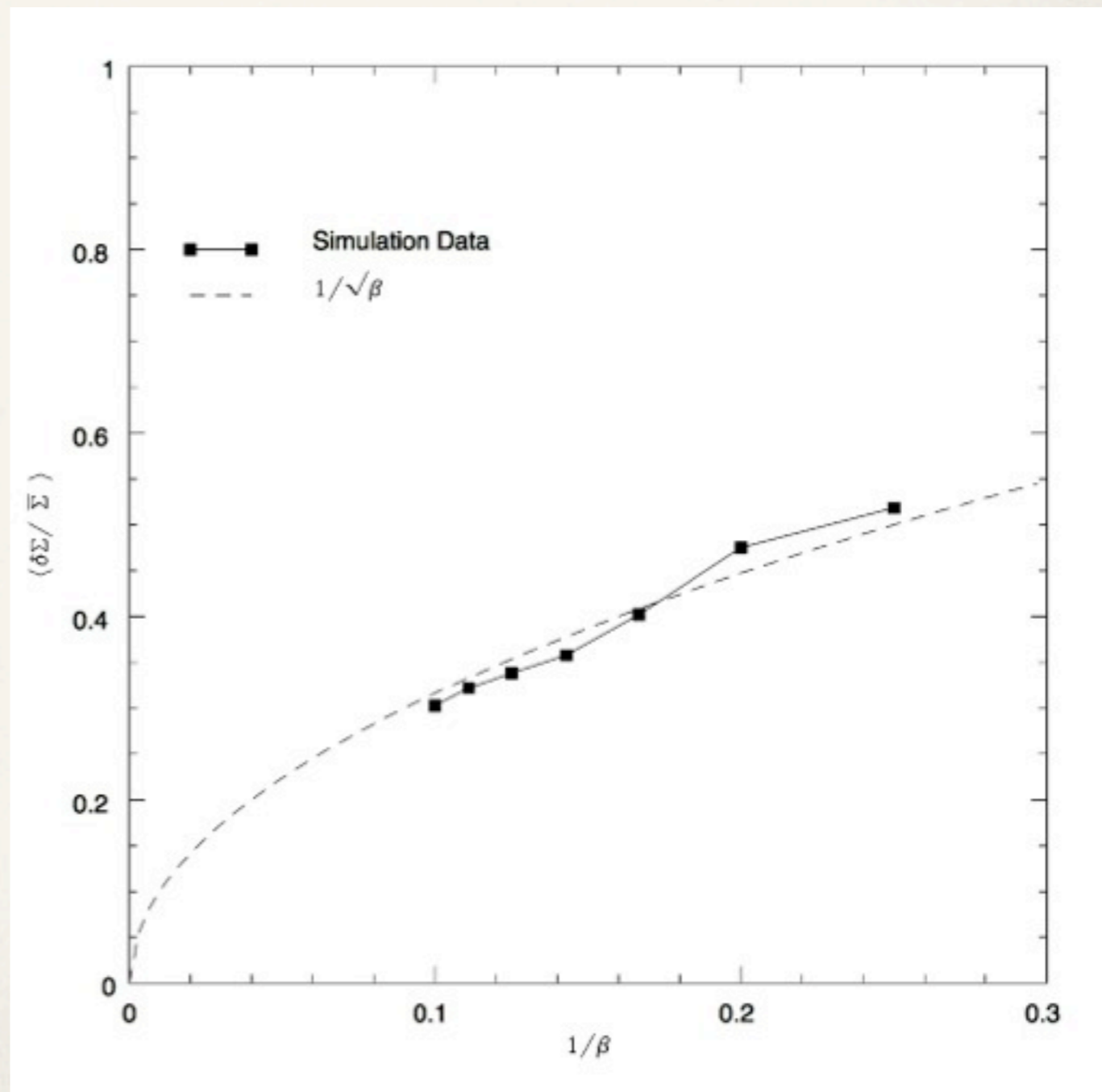
Cossins, Lodato & Clarke 2009

- * Self-regulation is established through thermal saturation of the spiral waves.
- * Amplitude of density perturbation must be related to cooling rate

- * We find that:

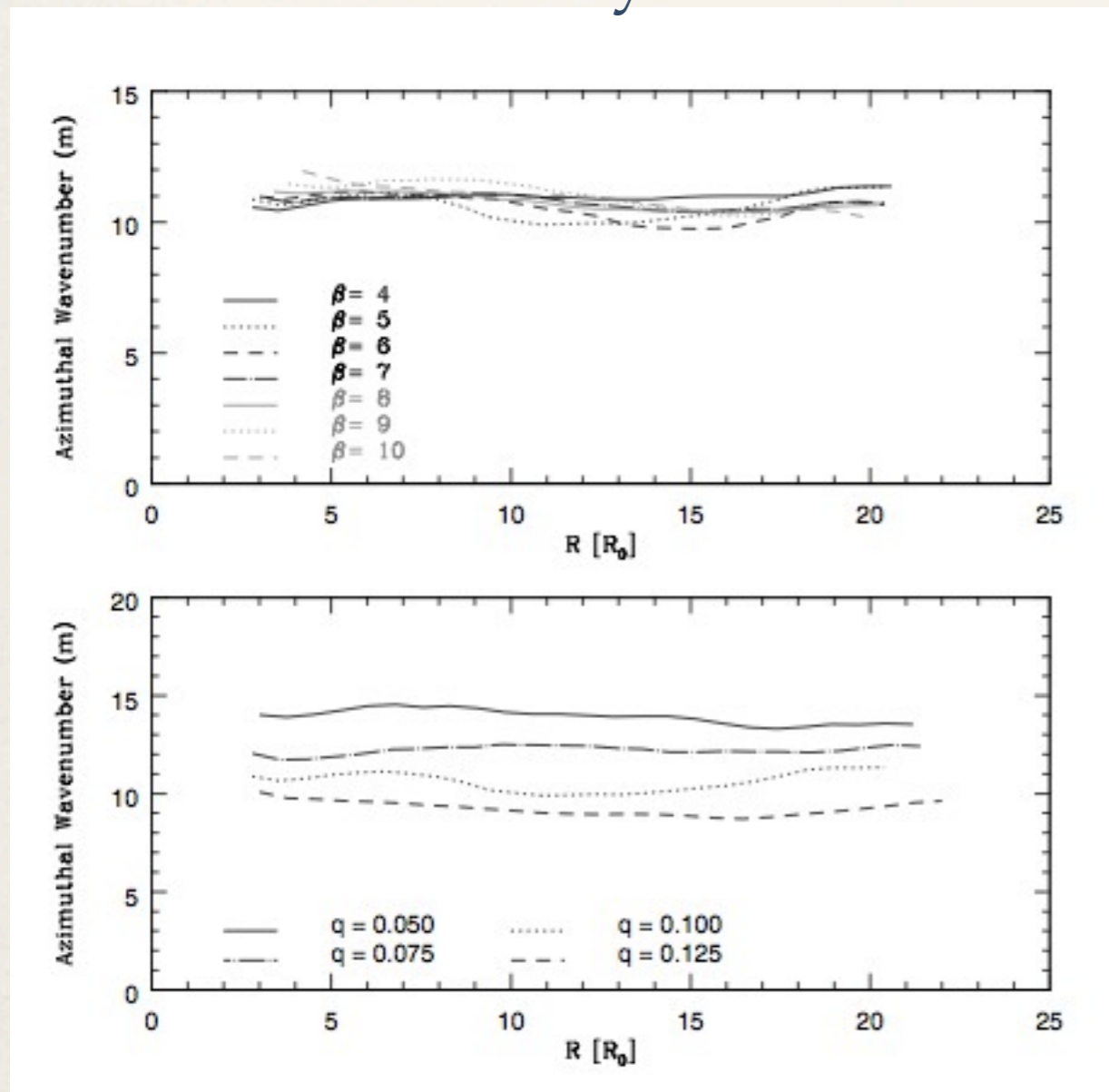
$$\frac{\Delta\Sigma}{\Sigma} \approx \frac{1}{\sqrt{\beta}}$$

- * Natural if consider that energy content of waves is proportional to the square of the perturbed fields

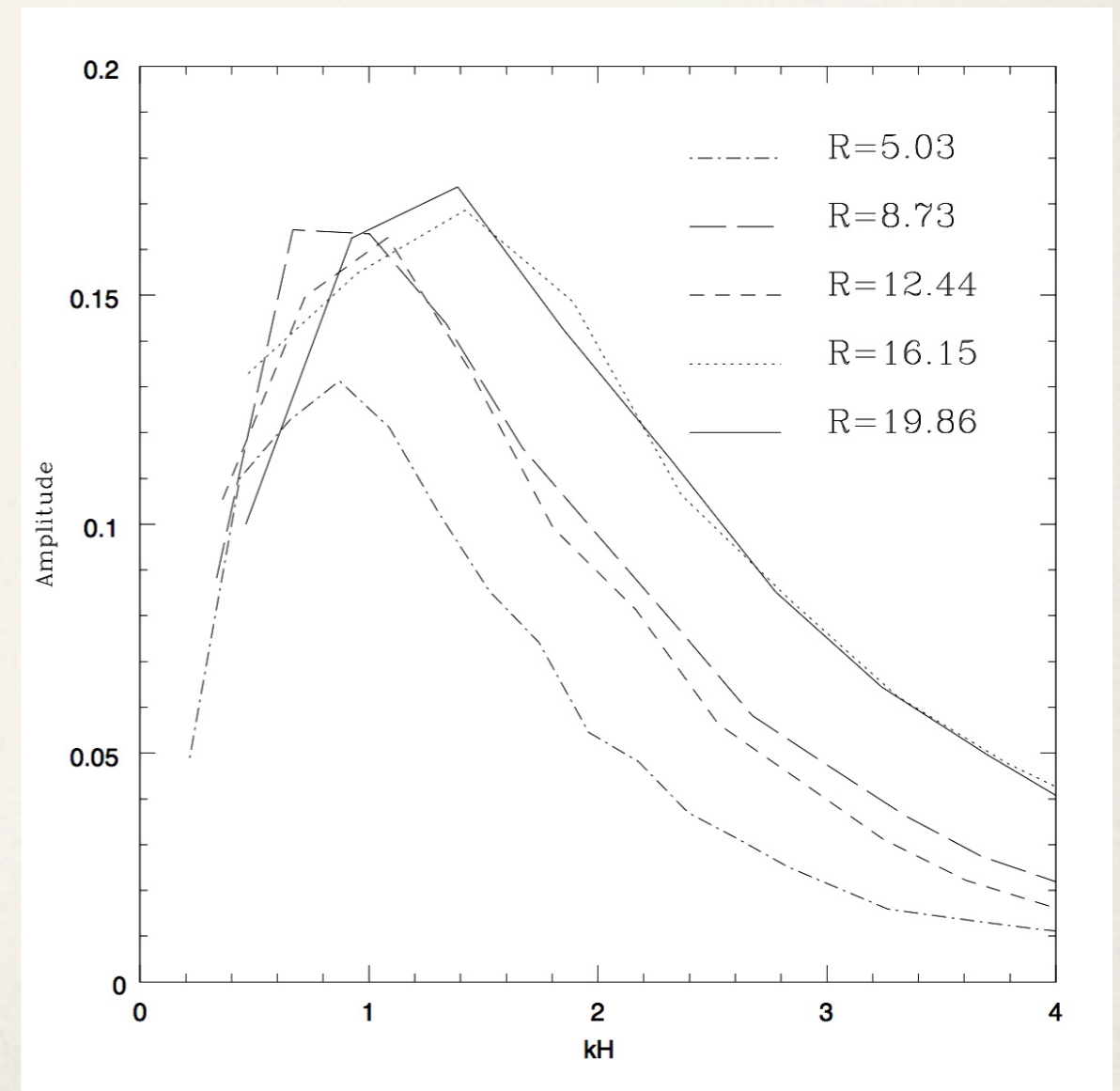


Spectrum of excited modes

Azimuthal structure: massive discs characterized by small m



Radial structure: at all radii, k peaks at roughly $\sim 1/H$

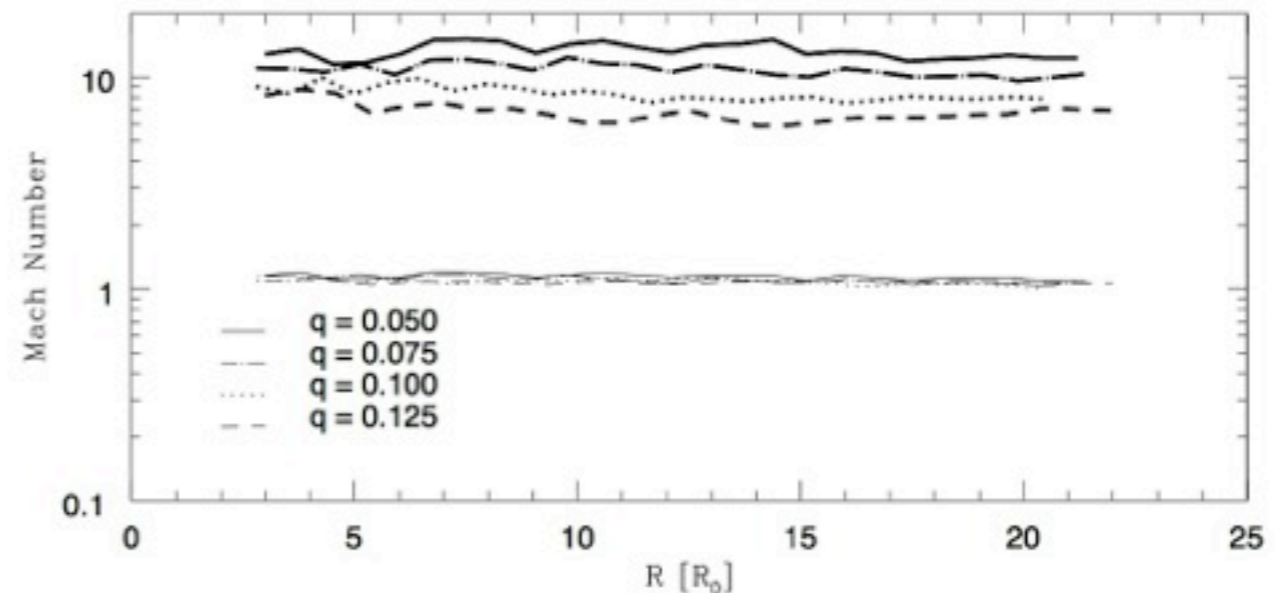
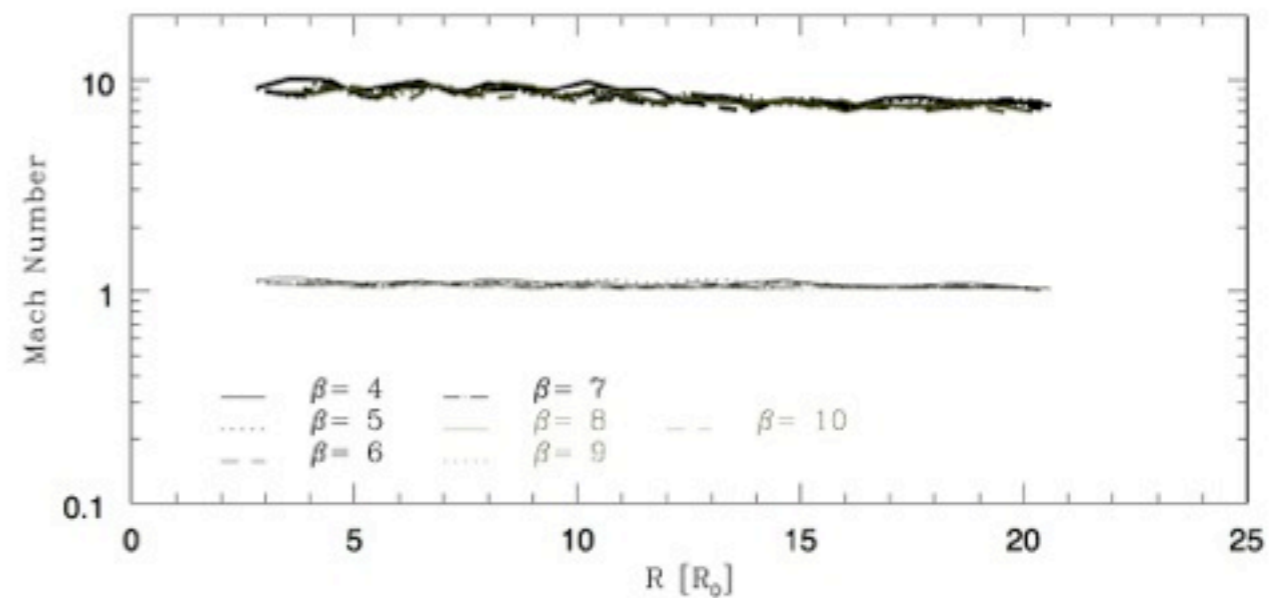


Sonic condition for spiral waves

Cossins, Lodato & Clarke 2009

- ❖ We have computed the pattern speed of the underlying spiral structure and its Mach number
- ❖ The Doppler-shifted Mach number is very close to unity, independently on radius, cooling rate, and disc mass.
- ❖ Density jump for almost sonic shocks also directly leads to

$$\frac{\Delta\Sigma}{\Sigma} \approx \frac{1}{\sqrt{\beta}}$$



Local vs global behaviour

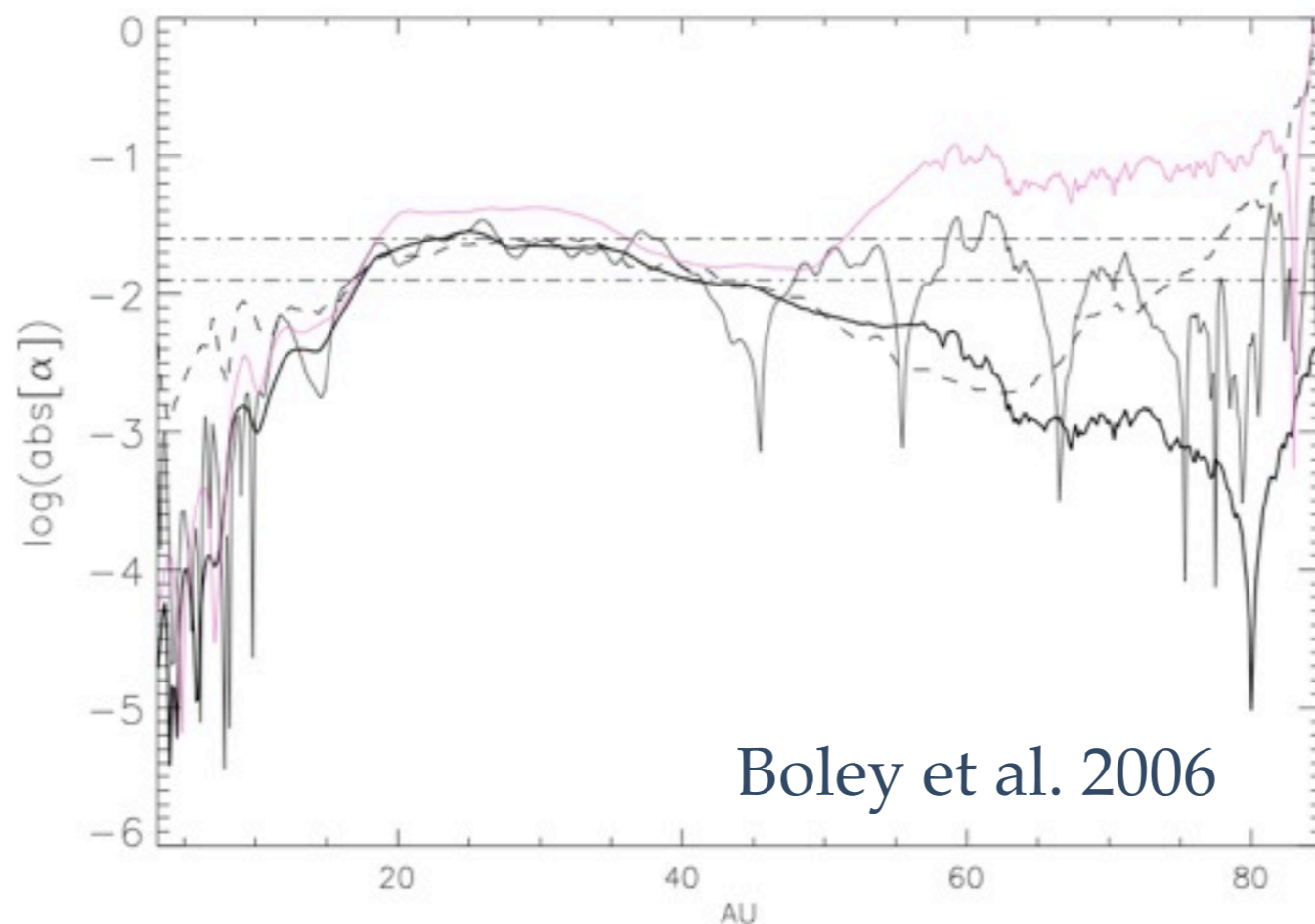
Lodato & Rice 2004

- * Can the evolution of self-gravitating discs be described within the standard, local, α -like prescription?
- * Can compute gravitational + Reynolds stresses directly from simulations and compare with expectations from standard α -theory (LR04, see also Boley et al. 2006)
- * **The disc adjusts so as to deliver the viscosity needed to stay in thermal equilibrium**

Local vs global behaviour

Lodato & Rice 2004

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- ❖ Can compute gravitational + Reynolds stresses directly from simulations and compare with α prescription (Lodato & Rice 2004, also Boley et al. 2006)
- ❖ **The disc ad equilibrium**



thermal

Local vs global behaviour

Cossins, Lodato & Clarke 2009

- * Can the evolution of self-gravitating discs be described within the standard, local, α -like prescription?
- * Described in detail by Balbus & Papaloizou (1999), recently discussed extensively by Cossins et al (2009)

- * Relation between energy and angular momentum densities in a density wave

$$\mathcal{E} = \Omega_p \mathcal{L} \quad \longrightarrow \quad \dot{\mathcal{E}} = \Omega_p \dot{\mathcal{L}}$$

- * Relation between power and stress due to local (viscous) processes

$$\dot{\mathcal{E}}_\nu = \Omega \dot{\mathcal{L}}_\nu$$

- * If density waves dissipate far from co-rotation, behaviour is non-local

Local vs global behaviour

Cossins, Lodato & Clarke 2009

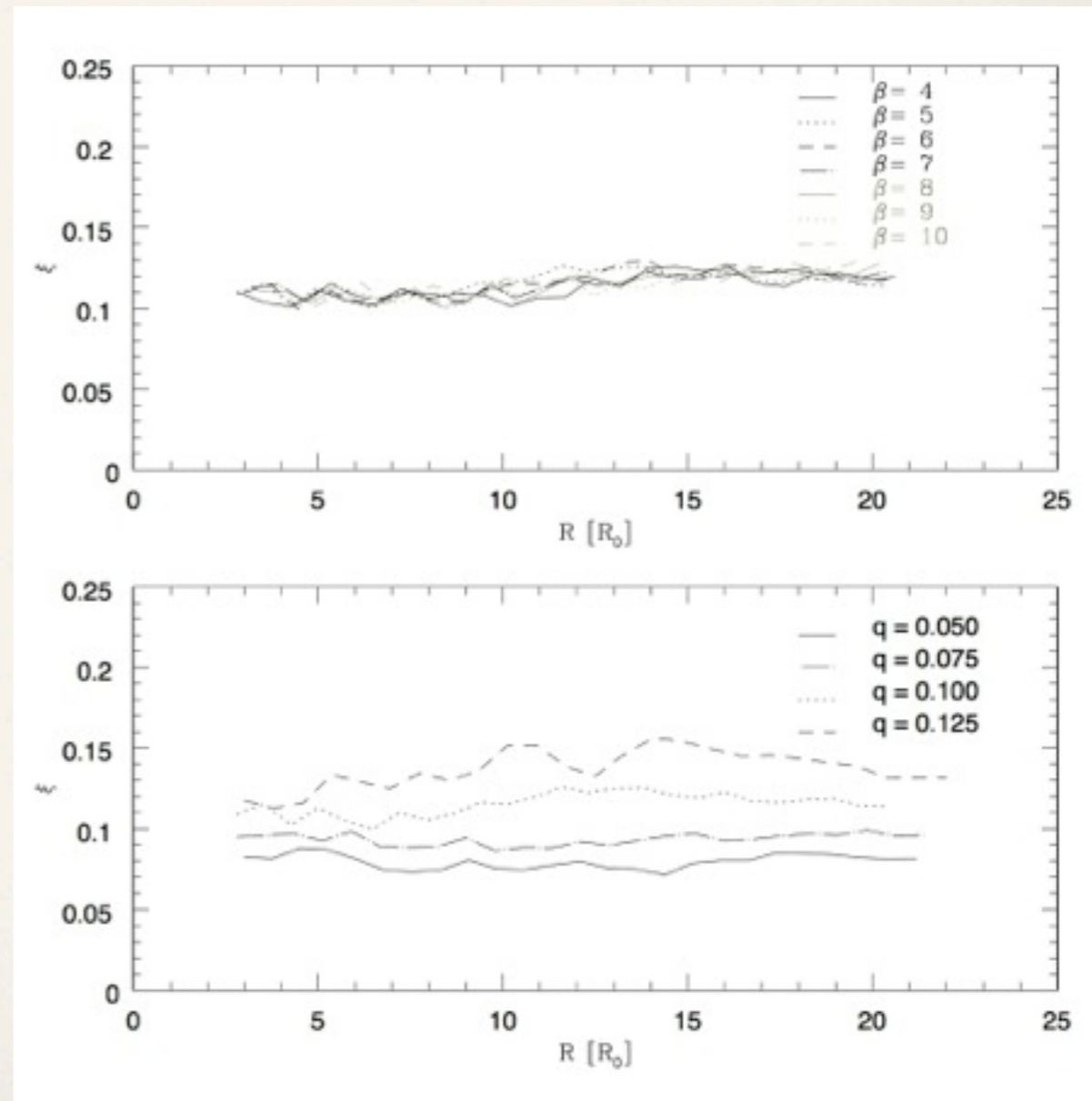
- * Degree of non-locality can be measured by

$$\xi = \left| \frac{\Omega - \Omega_p}{\Omega} \right|$$

- * Sonic condition for wave dissipation also tells us something about this:

$$\xi \approx \frac{c_s}{v_\phi} = \frac{H}{R}$$

- * To the extent that the disc is thin ($H \ll R$), global behaviour should be negligible
- * Possible to construct local, viscous models of disc evolution (Clarke 2009, Rafikov 2009)



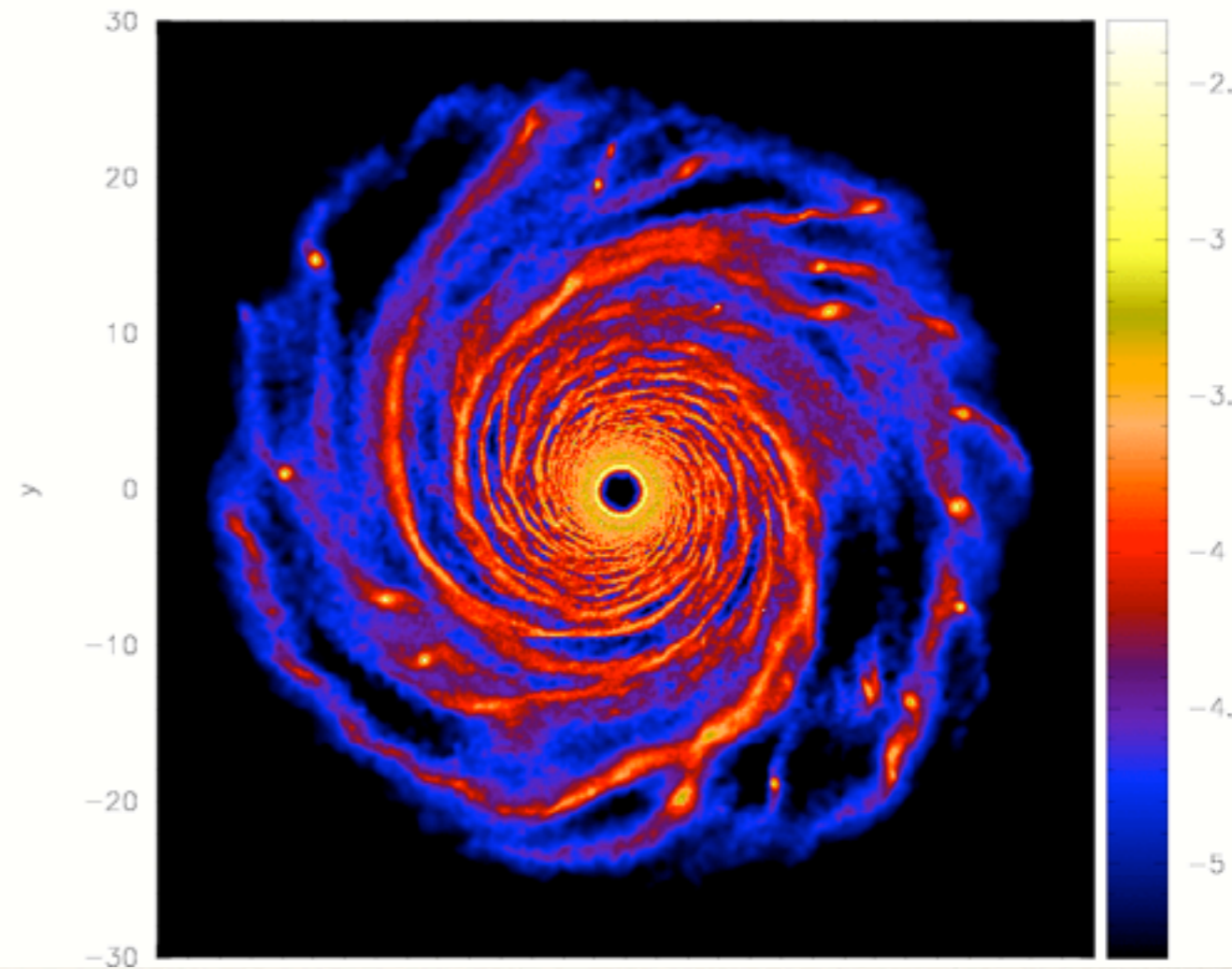
Convergence of numerical results

- ❖ It is well known that for short cooling times the disc is subject to fragmentation
- ❖ Meru and Bate (2011) show that such simulations are not converged
- ❖ As resolution increases fragmentation appears to become effective for longer cooling times

Simulation by Peter Cossins
 $\beta = 4$

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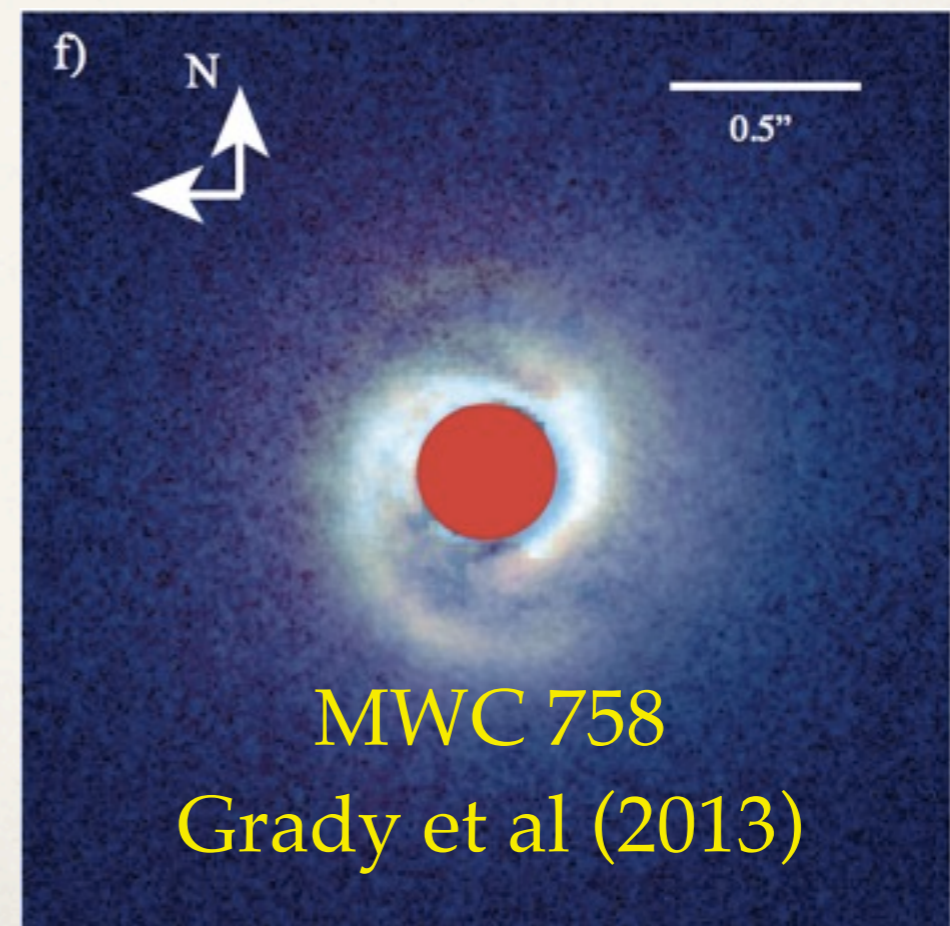
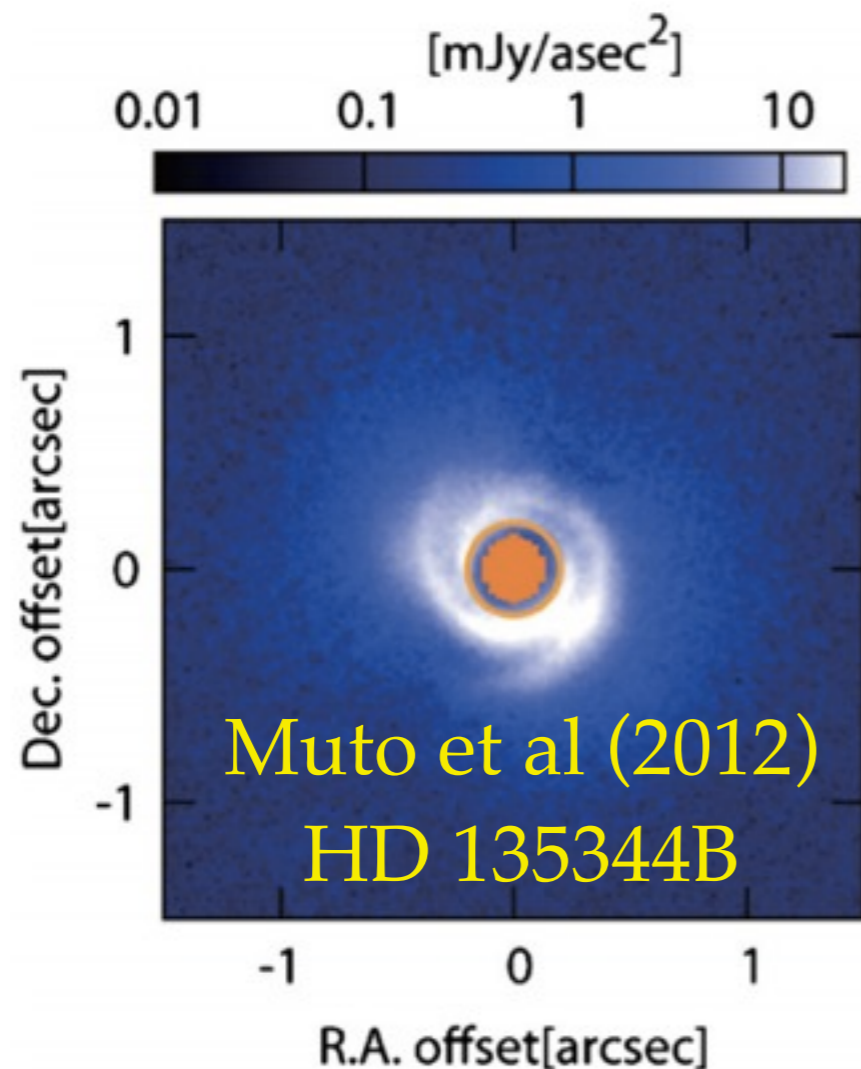
- ❖ Non-convergence most likely due to the effects of artificial viscosity in SPH simulations: additional viscosity (or any heating terms) weakens the gravitational instability (Lodato & Clarke 2011)
- ❖ Meru and Bate (2012): result converge at extremely high resolution. Fragmentation for $\beta < 20$
- ❖ Paardekooper (2012): non-convergence observed in 2D grid-based simulations. Fragmentation seen as a stochastic process
 - ❖ Essential to compute the likelihood of fragmentation in realistic protostellar discs (Hopkins and Christiansen 2013)

Convergence of numerical results

- ❖ What about convergence of results in the non-fragmenting limit?
- ❖ Is the stress and effective alpha converged?
- ❖ Michael et al (2012) (the Indiana group): convergence of measured stress observed in grid-based simulations
 - ❖ At low resolution alpha appears to be overestimated (more power in large scale structures ---> potentially more global transport)

Observing density waves in protostellar discs

- ❖ Several discs with spiral structures observed in scattered light
- ❖ Most of these are very evolved systems (transitional discs): most likely the origin of the spiral is not due to self-gravity

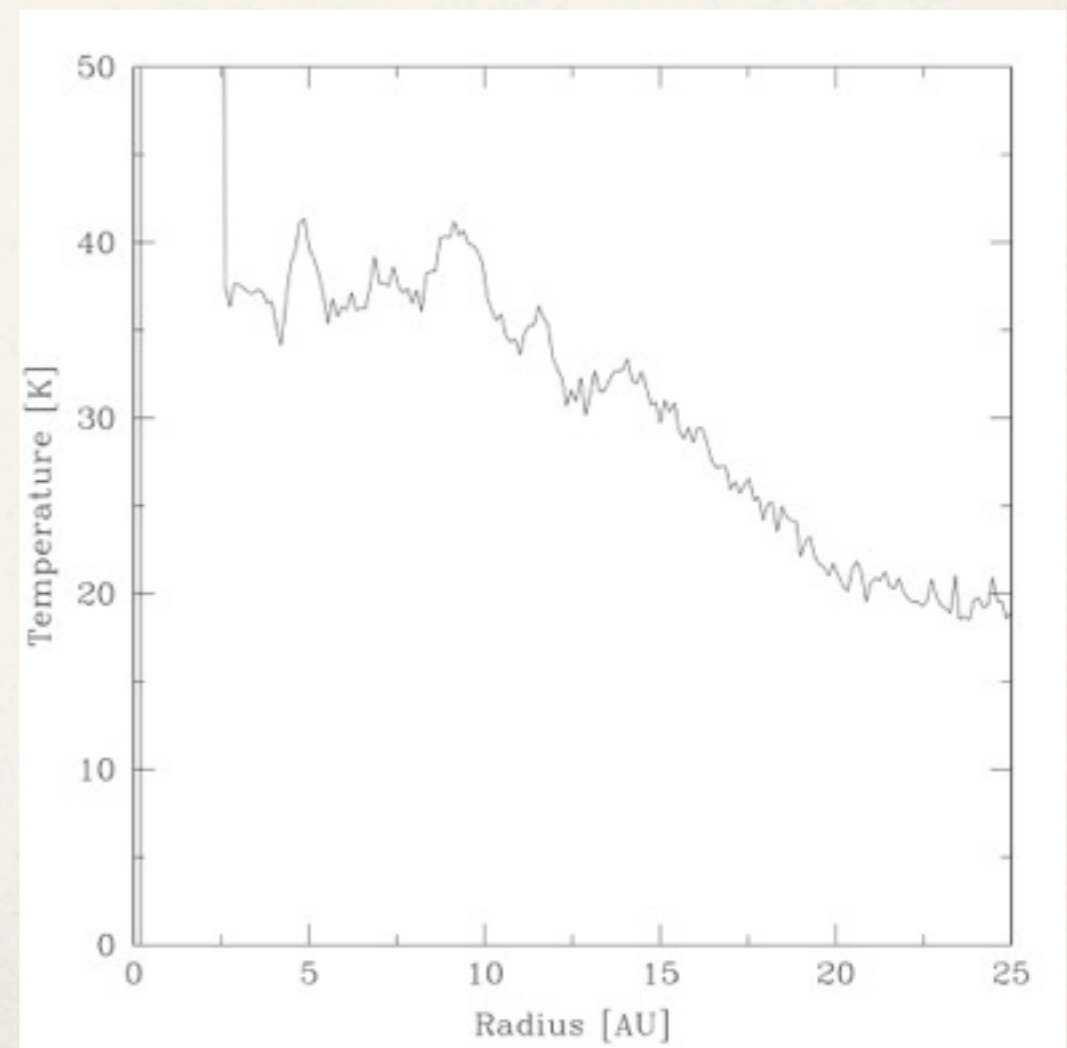
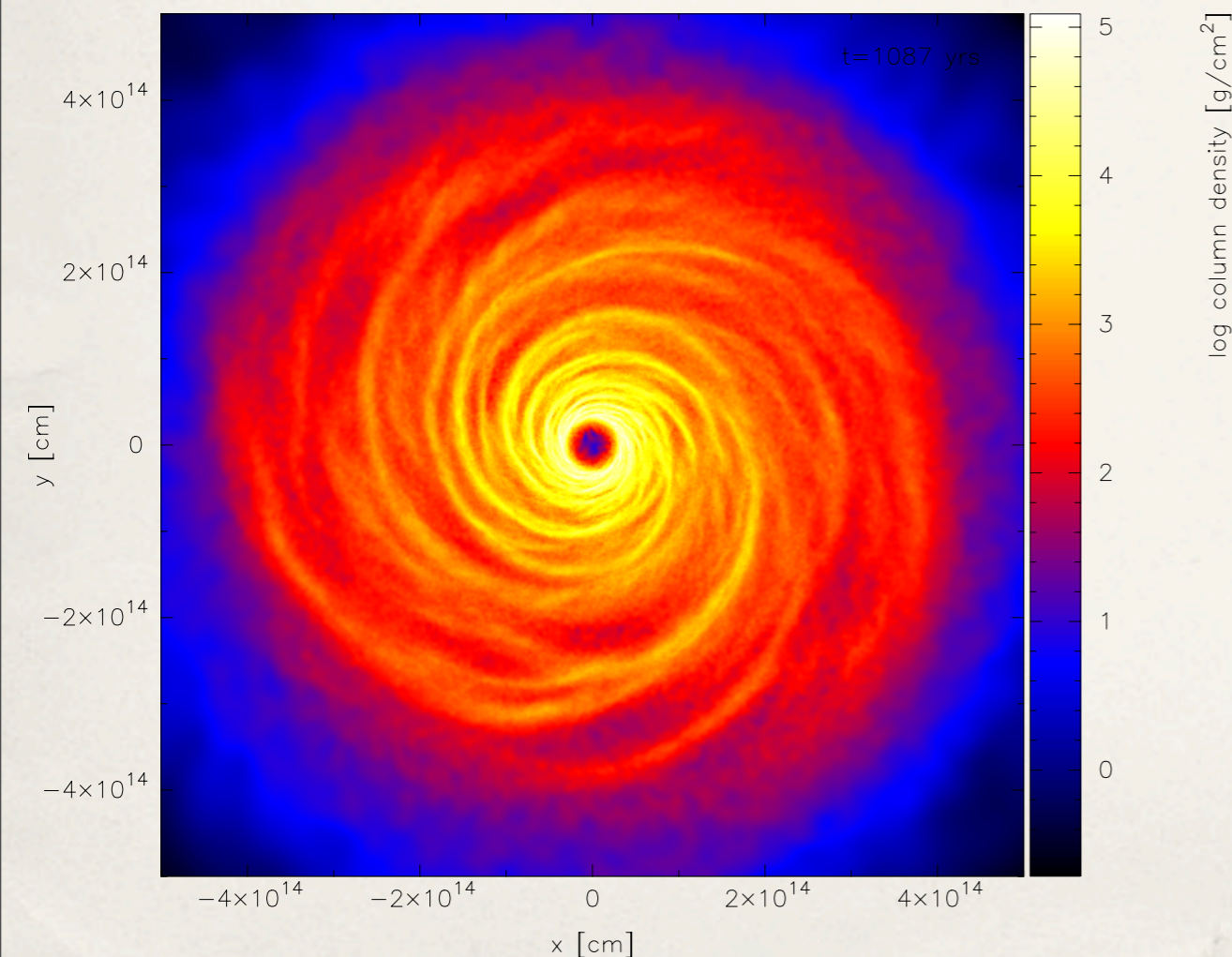


Self-gravitating discs with ALMA

Cossins, Lodato & Testi (2010)



- ❖ Start from one of our simulations, e.g. $M^* = 1M_{Sun}$, $M_{disc} = 0.2M^*$
- ❖ “Place” disc at 140pc (in Taurus) or at 50 pc (distance to TW Hya)
- ❖ Assume a “standard” opacity law

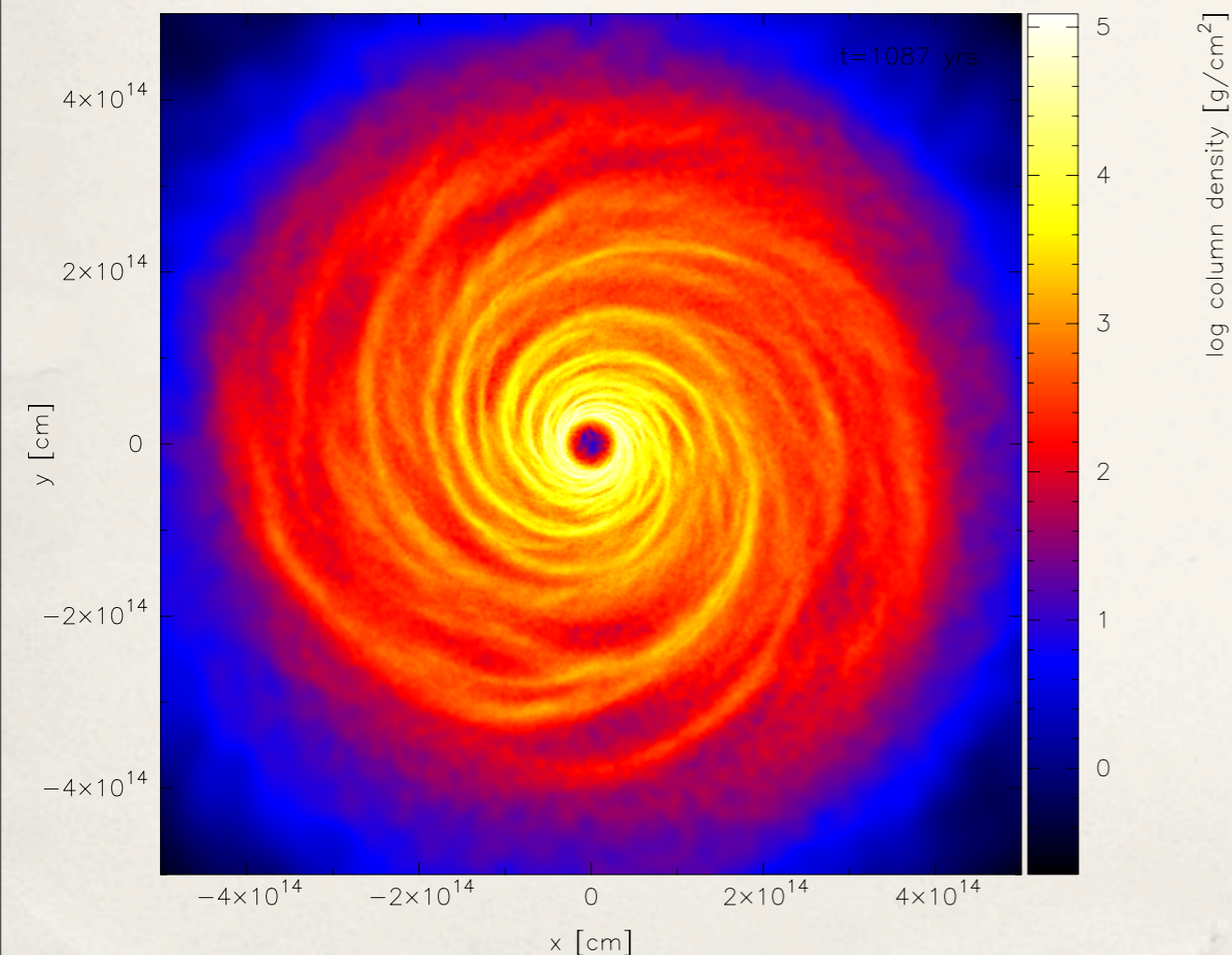


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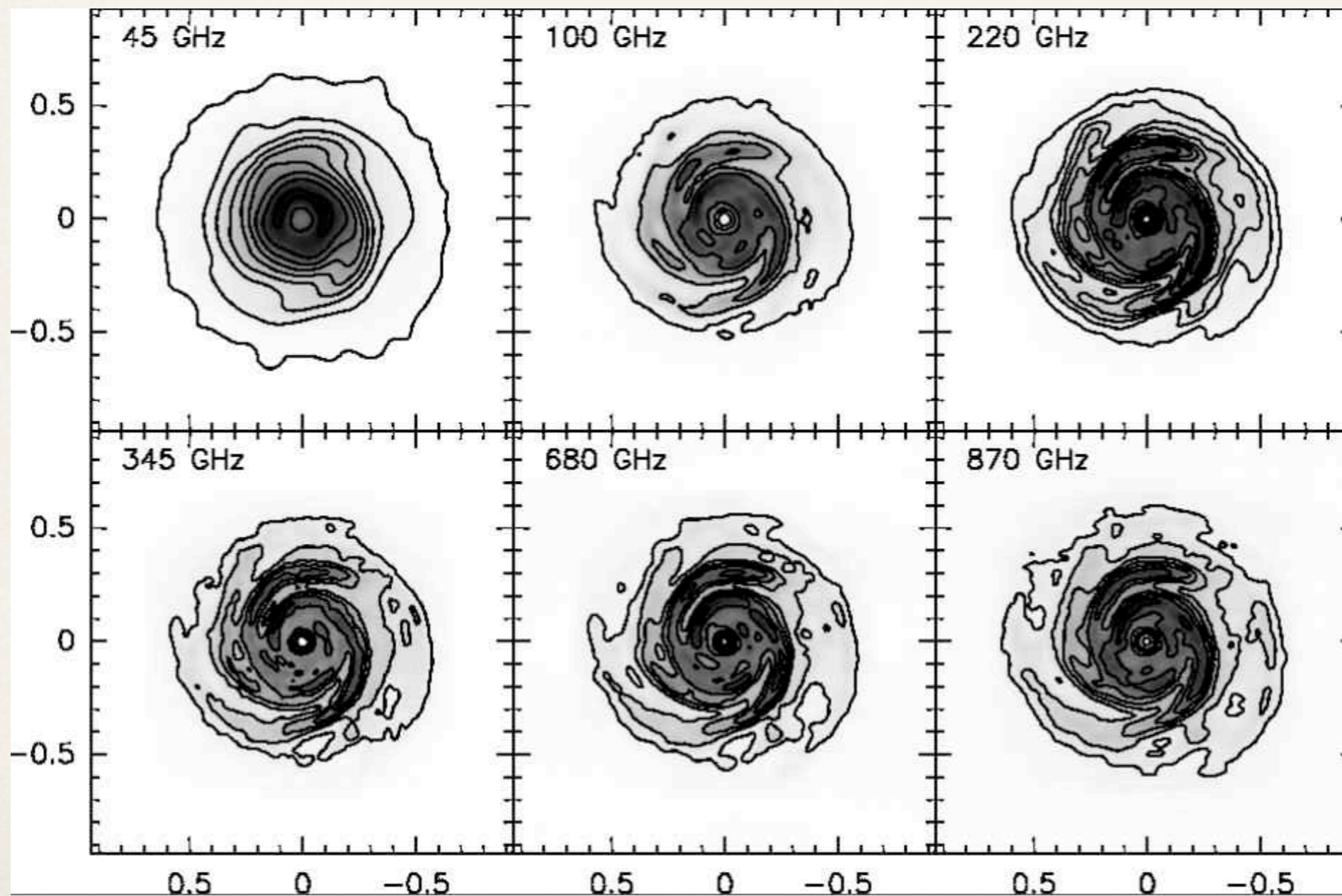


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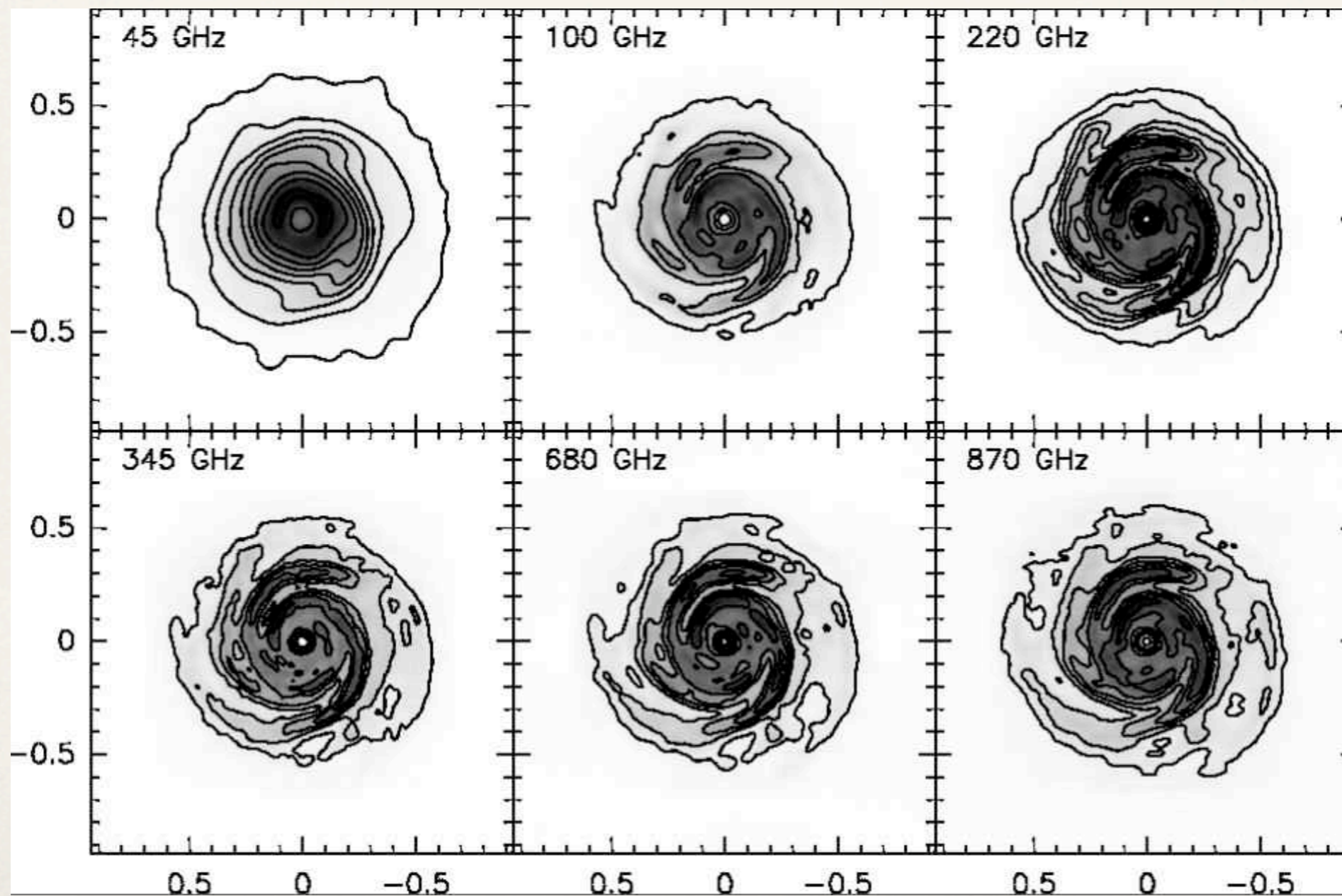


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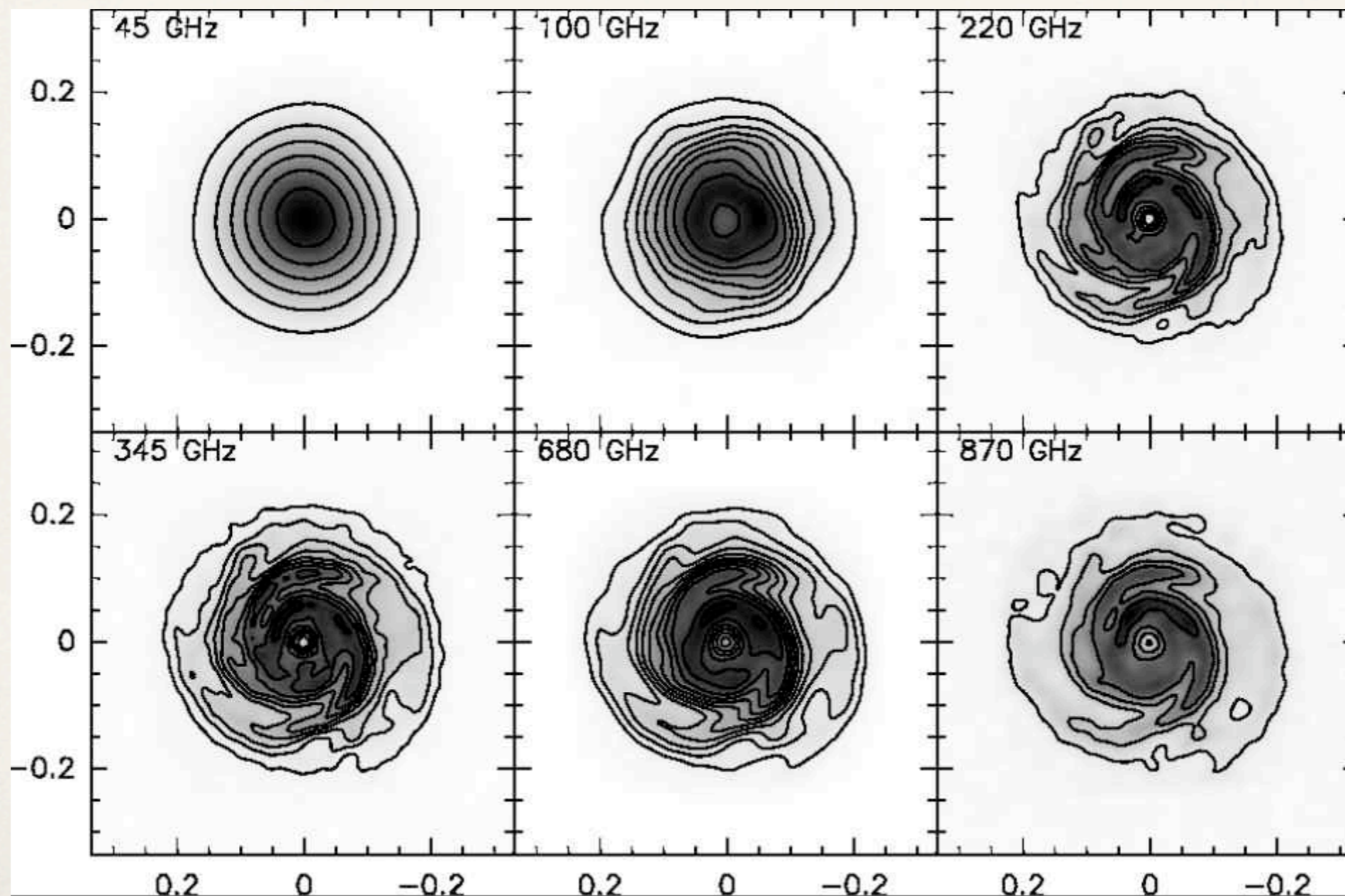


Simulated ALMA images at 50 pc



Self-gravitating discs with ALMA

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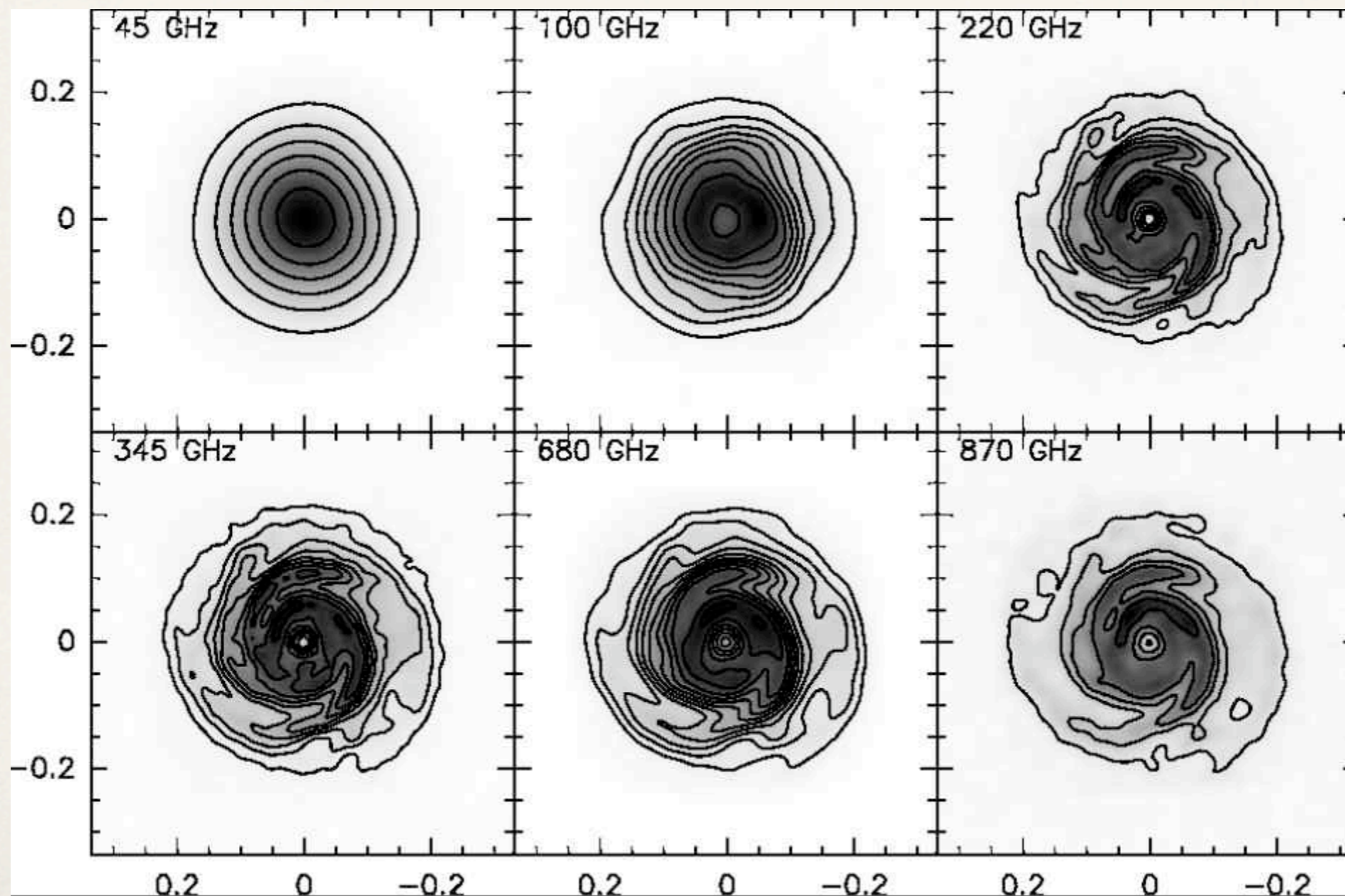


Self-gravitating discs with ALMA

Cossins, Lodato & Testi (2010)



Simulated ALMA images at 140 pc



Conclusions

- ❖ Young protostellar discs are likely to be gravitationally unstable
- ❖ Self-regulated evolution of GI leads to sustained angular momentum transport for ~ 1 Myr, bringing the disc into the T Tauri phase
- ❖ Density waves dissipate when they become sonic
- ❖ Induced transport is local IF disc is sufficiently thin
- ❖ GI could lead to fragmentation: exact fragmentation conditions unfortunately strongly affected by numerical resolution
- ❖ ALMA could be able to detect such discs very soon!